

linearize equations and we get

• Mass conservation: $\eta_t + gH u_x = 0$

• Momentum conservation: $H u_t + gH \eta_x = 0$

$$\begin{cases} \eta_t + gH u_x = 0 \\ u_t + g\eta_x = 0 \end{cases}$$

\Downarrow

$$\eta_{tt} - \underbrace{gH}_{c^2} \eta_{xx} = 0$$

wave equation.

Speed of Tides:

$$c = \sqrt{gH} \approx \sqrt{10 \times 1000} = 200 \text{ m s}^{-1}$$

\swarrow 3m deep $\quad \quad \quad \approx 720 \text{ km h}^{-1}$

Back to non-linear equations. Deriving momentum conservation

$$h u_t + h u u_x + g h h_x = 0 \implies u_t + u u_x + g h_x = 0$$

$$\underbrace{+(h u)_x + h_t}_{=0 \text{ By mass conservation}}$$

\swarrow we forgot about shocks for a while

$$\begin{pmatrix} h \\ u \end{pmatrix}_t + \underbrace{\begin{pmatrix} u & h \\ g & u \end{pmatrix}}_A \begin{pmatrix} h \\ u \end{pmatrix}_x = 0$$

In the linear case the matrix is $\begin{pmatrix} 0 & H \\ g & 0 \end{pmatrix}$.

we're going to pick vector (l_1, l_2) and consider

$$(l_1, l_2) \begin{pmatrix} h \\ u \end{pmatrix}_x + (l_1, l_2) A \begin{pmatrix} h \\ u \end{pmatrix}_x = 0$$

we will choose (l_1, l_2) as left eigenvector of A .

Module - 3

Complex potential $\Phi(z) = \phi(x,y) + i\psi(x,y)$.


velocity potential \rightarrow stream function

$$(u, v) = \nabla \phi$$

$$\frac{\partial \Phi}{\partial z} = \frac{\partial \Phi}{\partial x} = -i \frac{\partial \Phi}{\partial y} = \phi_x + i\psi_x = -i(\phi_y + i\psi_y) = u - iv$$

$$(-v, u) = \nabla \psi$$

$$\frac{d\Phi}{dz} = u - iv = \frac{dx}{dt} - i \frac{dy}{dt} = \frac{d\bar{z}}{dt}$$

Ex1:  $\Phi = Uz$
complex potential



Milne Thompson circle theorem: Let $f(z)$ analytic in $|z| < R$. Then in the presence of the cylinder the relative potential is

$$\Phi(z) = f(z) + \overline{f\left(\frac{R^2}{\bar{z}}\right)} = \frac{R^2}{z}$$

cylinder

Essentially we have one condition on the boundary of the problem.

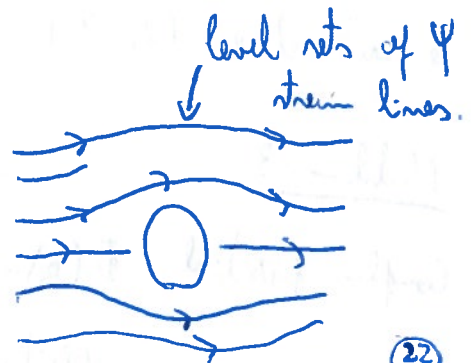
$$\Phi(R e^{i\theta}) = 2 \operatorname{Re} f(z) \in \mathbb{R} \Rightarrow \psi = 0$$

Ex: $\psi(x, y) = \text{constant}$ are tangent to the velocity vector

$$z^* = R^2/\bar{z} \text{ image point of } z$$

Using Milne Thompson Th in ex2,

$$\Phi(z) = U\left(z + \frac{R^2}{z}\right) \Rightarrow \psi(x, y) = Uy\left(1 - \frac{R^2}{x^2 + y^2}\right)$$



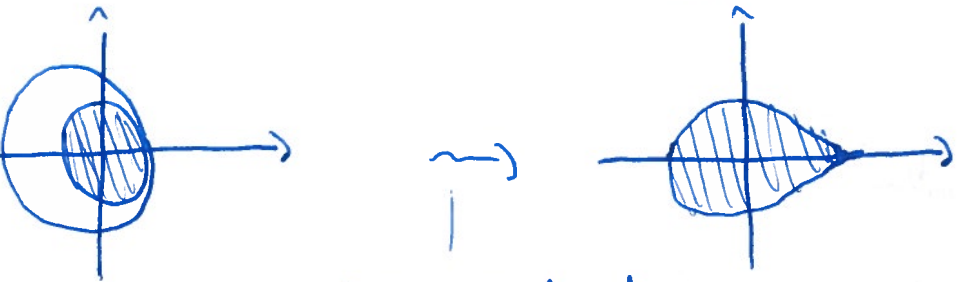
$z = \frac{1}{2}(w + \frac{\lambda^2}{w}) \rightarrow$ Joukowski transformation

Maps a circle to an ellipse.

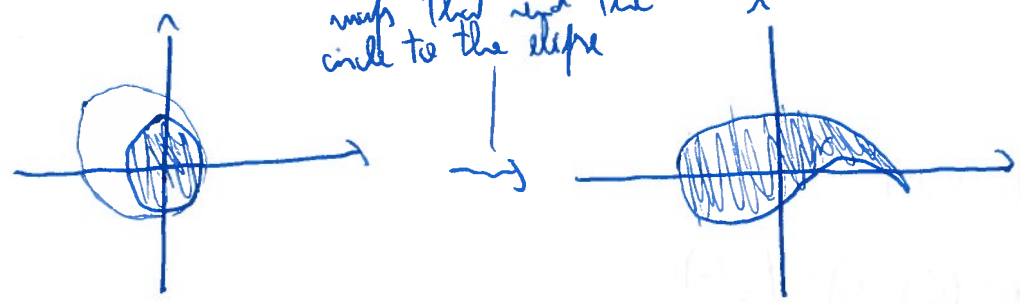
By choosing λ^2 accordingly you can make the ellipse collapse.

Ref: Ablowitz & Fokas, Cambridge University Press.

Airplane



By the conformal maps that send the circle to the ellipse



Ex 3: $\Phi(z) = -i\alpha \log z$

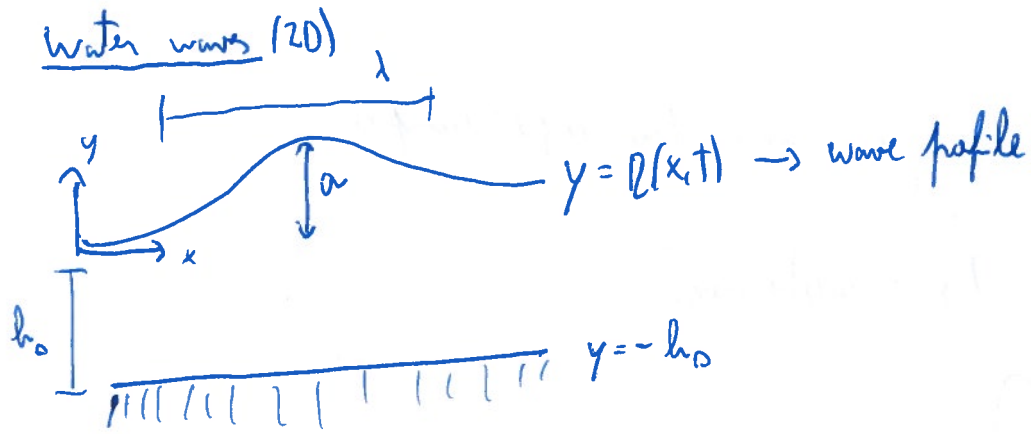


Wing of airplane



Nochlin 4

Water waves (2D)



Assumptions: inviscid, incompressible, irrotational, impermeable soil, etc.

Eq:

$$\begin{cases} \Delta \phi = 0 \\ \frac{d\phi}{dn} = 0 \text{ in the bottom} \\ \phi_t + \frac{1}{2}(\phi_x^2 + \phi_y^2) + g\eta = 0 \\ \eta_t + \phi_x \eta_x - \phi_y = 0 \end{cases}$$

$\phi = \phi(x, y, t)$ vel potential, $\nabla \phi = (u, w)$

Initial conditions:

$$\begin{cases} \phi(x, \eta(x, 0), 0) = \phi_0(x) \\ \eta(x, 0) = \eta_0(x) \end{cases}$$

Nonlinear Potential Theory

Whitham - Linear and Nonlinear Waves.

New variables: $x = \lambda \bar{x}$, $y = h_0 \bar{y}$, $t = \frac{\lambda}{c_0} \bar{t}$, $\eta = a \bar{\eta}$, $\phi = \frac{g \lambda a}{c_0} \bar{\phi}$

non-dimensional.

Dimensionless equations (drop the ~)

1) $\beta \phi_{xx} + \phi_{yy} = 0$ ← "Laplace" eq.

2) $\Omega_t + d \phi_x \Omega_x - \frac{1}{\beta} \phi_y = 0$ (*) at $y = d \eta$
 3) $\Omega + \phi_x + \frac{1}{2} d \phi_x^2 + \frac{1}{2} \frac{d}{\beta} \phi_y^2 = 0$

14) $\frac{d\phi}{dy} = 0, y = -1$ at $y = \pm d \eta$

where $d = \frac{a}{b_0}, \beta = \frac{b_0^2}{\lambda^2}$.

If $d \rightarrow 0$ we get the linearized problem.

d - Non-linearity parameter.

β - dispersive parameter.

Regime $d = O(\epsilon), \beta = O(\epsilon), \epsilon \ll 1$.

weakly nonlinear, weakly dispersive regime.

Ansatz: $\Omega, \phi(x, y, t) = \sum_m y^m f_m(x, t)$ ← Separation of variables.

Plug Ω, ϕ in "Laplace eq" (1) and impose Neumann condition (4)

we get

$$\phi(x, y, t) = \sum_{m=0}^{\infty} (-\beta)^m \frac{y^{2m}}{(2m)!} \partial_x^{2m} f(x, t) \quad \text{where}$$

$f = f_0$

Now we substitute ϕ in (2), (3)
 ↳ Non-linear.

Leading order:

$$\phi_x = f_x - \frac{\beta}{2} (1 + d\eta)^2 f_{xxx} + \underbrace{O(\beta^2)}_{O(\epsilon^2)}$$

$$\phi_y = -\beta (1 + d\eta) f_{xx} + \underbrace{O(\beta^2)}_{O(\epsilon^2)}$$

We only want to keep linear terms in d and β . I.e. truncate stuff beyond $O(d^2, d\beta, \beta^2)$.

Put this in (2) and (3) and truncate

$$\begin{cases} (2) & \eta_t + d f_x \eta_x + (1 + d\eta) f_{xx} - \frac{\beta}{6} f_{xxxx} = 0 \\ (3) & \eta_t f_t + \frac{\beta}{2} f_{xxt} + \frac{d}{2} f_x^2 = 0 \end{cases}$$

Burgers equations: η and $u \sim$ average of ϕ over y .

$$\frac{1}{Y} \int_0^Y \phi_x(x, y, t) dy \equiv U(x, t)$$

Table 4 -

$$\begin{cases} h_t + (h u)_x = 0 \\ u_t + u u_x + h_x = 0 \end{cases} \Leftrightarrow \begin{pmatrix} h \\ u \end{pmatrix}_t + \begin{pmatrix} u & h \\ 1 & u \end{pmatrix} \begin{pmatrix} h \\ u \end{pmatrix}_x = 0$$

In the linearized system:

$$\begin{cases} h_t + u_x = 0 \\ u_t + h_x = 0 \end{cases} \Rightarrow \begin{cases} (h+u)_t + \overbrace{(h+u)}^R_x = 0 \\ \underbrace{(h-u)}_L_t - (h-u)_x = 0 \end{cases}$$

leading order:

$$\phi_x = f_x - \frac{\beta}{2} (1 + \alpha \Omega)^2 f_{xxx} + \underbrace{O(\beta^2)}_{O(\epsilon^2)}$$

$$\phi_y = -\beta (1 + \alpha \Omega) f_{xx} + \underbrace{O(\beta^2)}_{O(\epsilon^2)}$$

we only want to keep linear terms in α and β . I.e. truncate stuff beyond $O(\alpha^2, \alpha\beta, \beta^2)$.

put this in (2) and (3) and truncate

$$\begin{cases} (2) & \Omega_t + \alpha f_x \Omega_x + (1 + \alpha \Omega) f_{xx} - \frac{\beta}{6} f_{xxxx} = 0 \\ (3) & \Omega + f_t + \frac{\beta}{2} f_{xxt} + \frac{\alpha}{2} f_x^2 = 0 \end{cases}$$

Boussinesq equations: Ω and $u \sim$ average of ϕ over y .

$$\frac{1}{Y} \int_0^Y \phi_x(x, y, t) dy \equiv U(x, t)$$

Table 4 -

$$\begin{cases} h_t + (h u)_x = 0 \\ u_t + u u_x + h_x = 0 \end{cases} \Leftrightarrow \begin{pmatrix} h \\ u \end{pmatrix}_t + \begin{pmatrix} u & h \\ 1 & u \end{pmatrix} \begin{pmatrix} h \\ u \end{pmatrix}_x = 0$$

In the linearized system:

$$\begin{cases} h_t + u x = 0 \\ u_t + h_x = 0 \end{cases} \Rightarrow \begin{cases} (h+u)_t + \overbrace{(h+u)}^R_x = 0 \\ \underbrace{(h-u)}_L_t - (h-u)_x = 0 \end{cases}$$