

The main idea of the **neural field models** in Mathematical Neuroscience is to **treat the cortical space as continuous**. Since the number of neurons and synapses is extremely high even in a small piece of cortex, this idea appears naturally as a first approximation to model the neural activity. This approach was first developed in 70's by Wilson and Cowan [5] and Amari [1]; it leads to integro-differential equations (or systems of them), which may be written in the form:

$$c \frac{\partial}{\partial t} V(\bar{x}, t) = I(\bar{x}, t) - V(\bar{x}, t) + \int_{\Omega} K(\|\bar{x} - \bar{y}\|_2) S(V(\bar{y}, t)) d\bar{y}, \quad (1)$$

$$t \in [0, T], \bar{x} \in \Omega \subset \mathbb{R}^2;$$

- ▶ $V(\bar{x}, t)$ - the membrane potential in point x at time t ;
- ▶ I - external sources of excitation;
- ▶ S - dependence between the firing rate of the neurons and their membrane potentials (sigmoidal or Heaviside function);
- ▶ $K(\|\bar{x} - \bar{y}\|_2)$ - connectivity between neurons at \bar{x} and \bar{y} .

Initial Condition: $V(\bar{x}, 0) = V_0(\bar{x})$, $\bar{x} \in \Omega$.

Numerical algorithms for the approximation of the Neural Field Equation in two dimensions have been proposed by Faye and Faugeras [3], Hutt and Rougier [4]. Here we propose a new numerical approach, based on the use of an **implicit second order scheme and Gaussian quadrature**.

Time Discretisation

We begin by rewriting equation (1) in the form

$$c \frac{\partial}{\partial t} V(\bar{x}, t) = I(\bar{x}, t) - V(\bar{x}, t) + \kappa(V(\bar{x}, t)) \quad (2)$$

$$t \in [0, T], \bar{x} \in \Omega \subset \mathbb{R}^2,$$

where

$$\kappa(V(\bar{x}, t)) = \int_{\Omega} K(\|\bar{x} - \bar{y}\|_2) S(V(\bar{y}, t)) d\bar{y}. \quad (3)$$

Let h_t be the stepsize in time. We define

$$t_i = ih_t, \quad i = 0, \dots, M, \quad T = h_t M.$$

Moreover, let $V_i(\bar{x}) = V(\bar{x}, t_i)$, $\forall \bar{x} \in \Omega$, $i = 0, \dots, M$. We shall approximate the partial derivative in time by the backward difference

$$\frac{\partial}{\partial t} V(\bar{x}, t_i) \approx \frac{3V_i(\bar{x}) - 4V_{i-1}(\bar{x}) + V_{i-2}(\bar{x})}{2h_t}, \quad (4)$$

which gives a **discretisation error of the order $O(h_t^2)$** , for sufficiently smooth V . By substituting (4) into (2) we obtain the implicit scheme

$$c \frac{3U_i - 4U_{i-1} + U_{i-2}}{2h_t} = I_i - U_i + \kappa(U_i), \quad i = 2, \dots, M, \quad (5)$$

where U_i approximates the solution of (2).

Space Discretisation

Assume that Ω is a rectangle: $\Omega = [-1, 1] \times [-1, 1]$. Introduce a uniform grid of points (x_i, x_j) , such that $x_i = -1 + ih$, $i = 0, \dots, n$, where h is the discretisation step in space. In each subinterval $[x_i, x_{i+1}]$ we introduce k Gaussian nodes: $x_{i,s} = x_i + \frac{h}{2}(1 + \xi_s)$, $i = 0, 1, \dots, n-1$, where ξ_s are the roots of the k -th degree Legendre polynomial, $s = 1, \dots, k$. Using a **Gaussian quadrature formula** to evaluate the integral, we obtain the finite-dimensional approximation of $\kappa(U)$. **This discretisation provides an accuracy order of $O(h^{2k})$** .

$$(\kappa^h(U^h))_{mu,lv} = \sum_{i=0}^{n_1} \sum_{j=0}^{n_2} \sum_{s=1}^k \sum_{t=1}^k \tilde{w}_s \tilde{w}_t \times K(\|(x_{mu,l}, x_{lv}) - (y_{is}, y_{jt})\|_2) S((U^h)_{is,jt}). \quad (6)$$

After discretization we obtain the following system of nonlinear equations:

$$U^h - \lambda \kappa^h(U^h) = f^h, \quad (7)$$

where $\kappa^h(U^h)$ is defined by (6) and $(f^h)_{is,jt} = f(x_{is}, x_{jt})$. To solve (7), which is a system of N^2 nonlinear equations, we use the **fixed point method**

Efficiency and Rank Reduction

In order to **improve the efficiency** of the numerical method, we apply the following technique, proposed in [6]. Assuming that the function V is sufficiently smooth, we can approximate it by an interpolating polynomial of a certain degree. **Our approach for reducing the matrices rank in our method consists in replacing the solution V_i by its m -th degree interpolating polynomial at the Chebyshev nodes in Ω** . If V_i is sufficiently smooth, this produces a very small error and yields a very significant reduction of computational cost. Actually, when computing the vector $\kappa^h(U^h)$ (see formula (6)) **we have only to compute $m^2 \ll N^2$ components**, one for each Chebyshev node.

Neural Field Equation with Delay

According to many authors (see for example [3]), **realistic models of neural fields must take into account that the propagation speed of neuronal interactions is finite**, which leads to NFE with delays of the form

$$c \frac{\partial}{\partial t} V(\bar{x}, t) = I(\bar{x}, t) - V(\bar{x}, t) + \int_{\Omega} K(\|\bar{x} - \bar{y}\|_2) S(V(\bar{y}, t - \tau(\bar{x}, \bar{y}))) d\bar{y}, \quad (8)$$

$t \in [0, T]$, $\bar{x} \in \Omega \subset \mathbb{R}^2$, where $\tau(\bar{x}, \bar{y}) > 0$ is a delay, depending on the spatial variables. Assuming that the electrical signals propagate with a constant speed v , uniform in space, we set $\tau(\bar{x}, \bar{y}) = \|\bar{x} - \bar{y}\|_2/v$. In the delay case, the initial condition has the form $V(\bar{x}, t) = V_0(\bar{x}, t)$, $\bar{x} \in \Omega$, $t \in [-\tau_{max}, 0]$, where $\tau_{max} = \max_{\bar{x}, \bar{y} \in \Omega} \tau(\bar{x}, \bar{y})$. **The numerical algorithm used to solve equation (8) is essentially the same as described in the previous sections**. The main difference results from the fact that when computing the integral on the right-hand side of (8) at instant t_i we must use the approximate solution at all instants t_{i-k} , $k = 1, \dots, K_{max}$, where K_{max} is the integer part of τ_{max}/h_t .

Numerical Example 1

In this example we consider a kind of neural field with spatially localized oscillations which occur in excitable neural media with short-range excitation and long-range inhibition (**mexican hat connectivity**), in the case of a **spatially localized input**. Such neural fields are known as **breathers** due to their oscillatory behaviour and their dynamics are analysed in [2]. A similar example was described and computed in [4]. In this case the connectivity kernel has the form:

$$K(r) = 20 \frac{\exp(-r)}{18\pi} - 14 \frac{\exp(-\frac{r}{3})}{18\pi}. \quad (9)$$

The firing rate function has the sigmoidal form

$$S(x) = \frac{2}{1 + \exp(-10000(x - 0.005))}$$

and the external input is given by

$$I(\bar{x}, t) = 5 \frac{\exp(-x^2/32 - y^2/32)}{32\pi}.$$

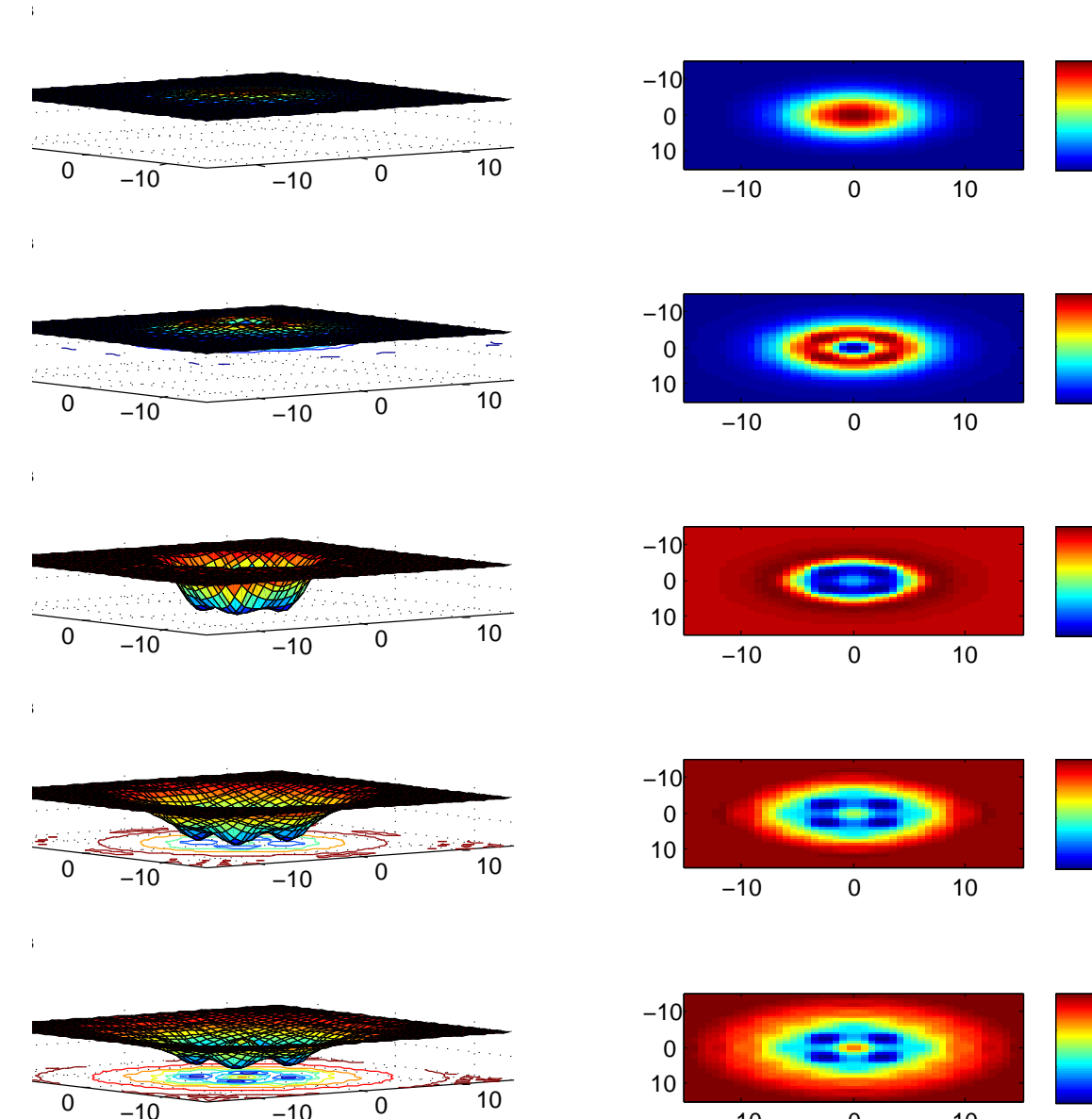


Figure 1. Plots of the solution in the case $v = 50$ at different moments of time: $t = 0.08, 0.16, 0.24, 0.32, 0.40$.

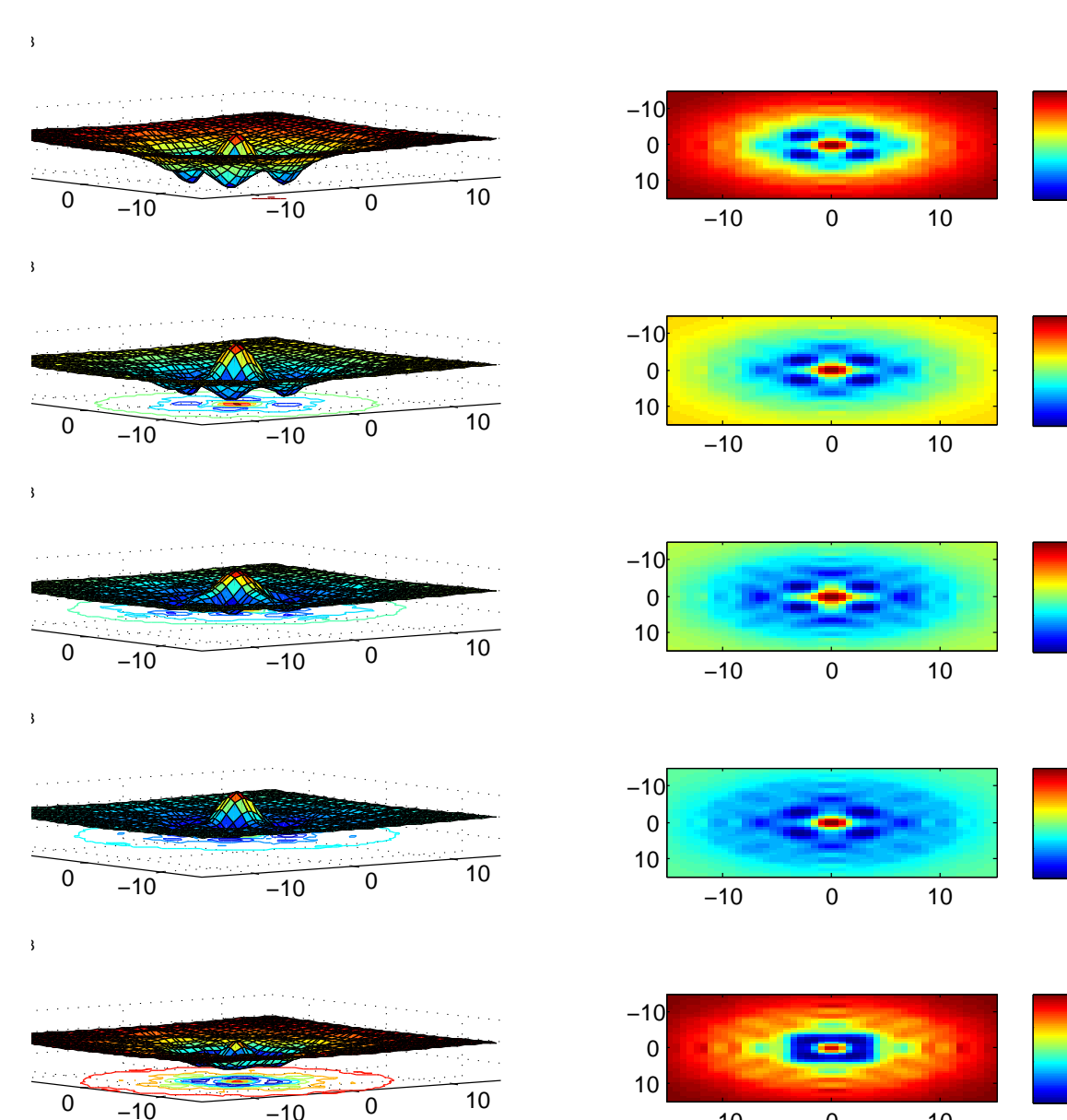


Figure 2. Plots of the solution in the case $v = 50$ at different moments of time: $t = 0.48, 0.56, 0.64, 0.72, 0.80$.

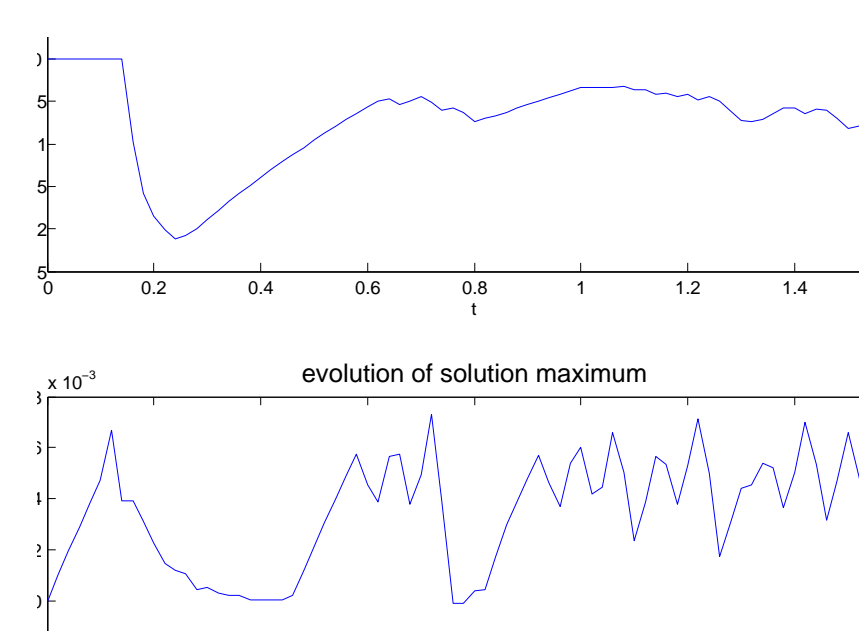


Figure 3. Evolution of the solution maximum and minimum for the breather with $v = 50$.

Numerical Example 2

In this example (described in [3]), the firing rate function has the form

$$S(x) = \frac{2}{1 + e^{-\mu x}}, \quad x \in \mathbb{R}.$$

where $\mu \in \mathbb{R}^+$, and the connectivity function is given by

$$K(r) = \frac{1}{\sqrt{2\pi\xi_1^2}} \exp\left(-\frac{r^2}{2\pi\xi_1^2}\right) - \frac{A}{\sqrt{2\pi\xi_2^2}} \exp\left(-\frac{r^2}{2\pi\xi_2^2}\right),$$

where $r = \|x - y\|_2 = \sqrt{(x_1 - y_1)^2 + (x_2 - y_2)^2}$ and $\xi_1, \xi_2, A \in \mathbb{R}^+$. We consider $c = 1$ and the external input $I(x, t)$ is 0. Our aim is to investigate **how the behaviour of the solutions depends on the equation parameters**. We set the initial condition $V(0, x) \equiv 0.01$ and check whether the solution tends or not to the trivial steady state. In Fig. 4 (left and right) a 3D-plot and a contour plot of the corresponding solution are displayed, respectively, for $t = 3$, in the case $A = 1$, $\xi_1 = 0.4$, $\xi_2 = 0.2$, $\mu = 10$, with no delay. The computations were carried on the time interval $[0, 3]$ with stepsize $h_t = 0.1$. The parameters of the space discretisation are $m = 12$, $N = 48$. Let x_1 be a point close to the center of Ω and x_2 be a point in the boundary of the domain. In Fig. 5 the graphs of $V(x_1, t)$ (left) and $V(x_2, t)$ (right), as functions of time, are displayed. The behaviour of the solution depends strongly on μ . For $\mu = 10$, for example, we see that after a certain time the solution becomes decreasing, both in x_1 and x_2 . But for $\mu = 15$, if t is sufficiently high, the solution increases in both points. This suggests that for some value of μ , between 10 and 15, there should be a **bifurcation** (the zero solution becomes unstable).

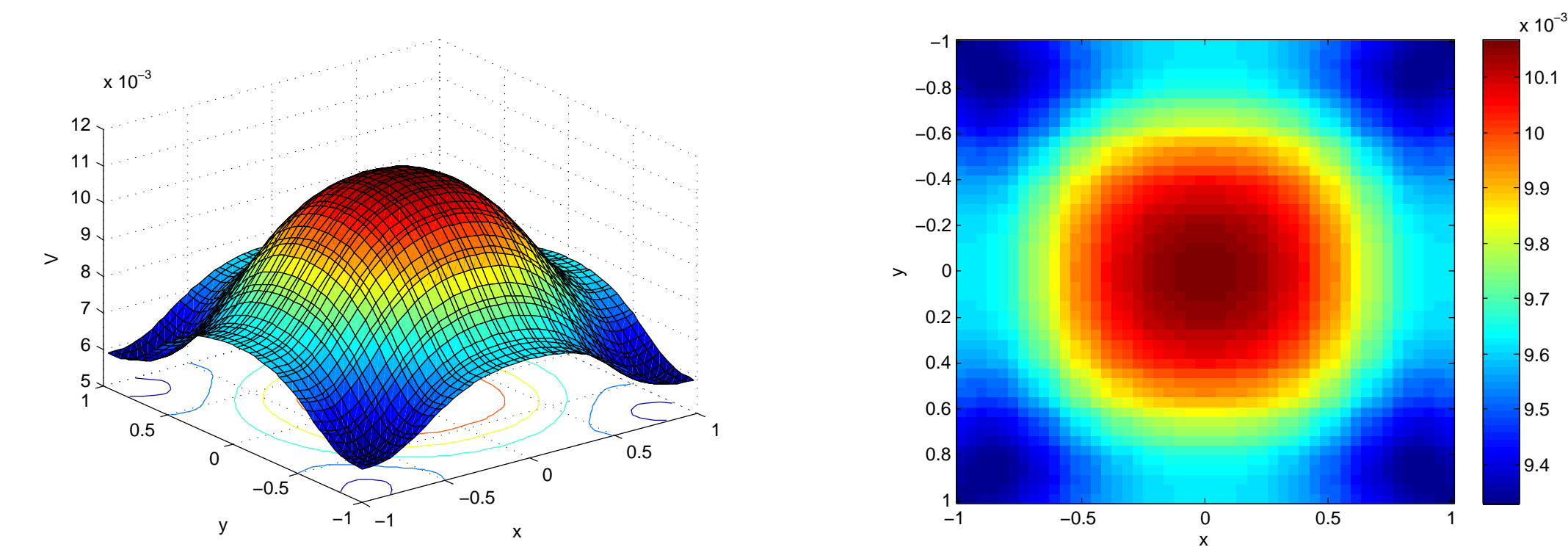


Figure 4. 3D-plot and contour plot of the solution

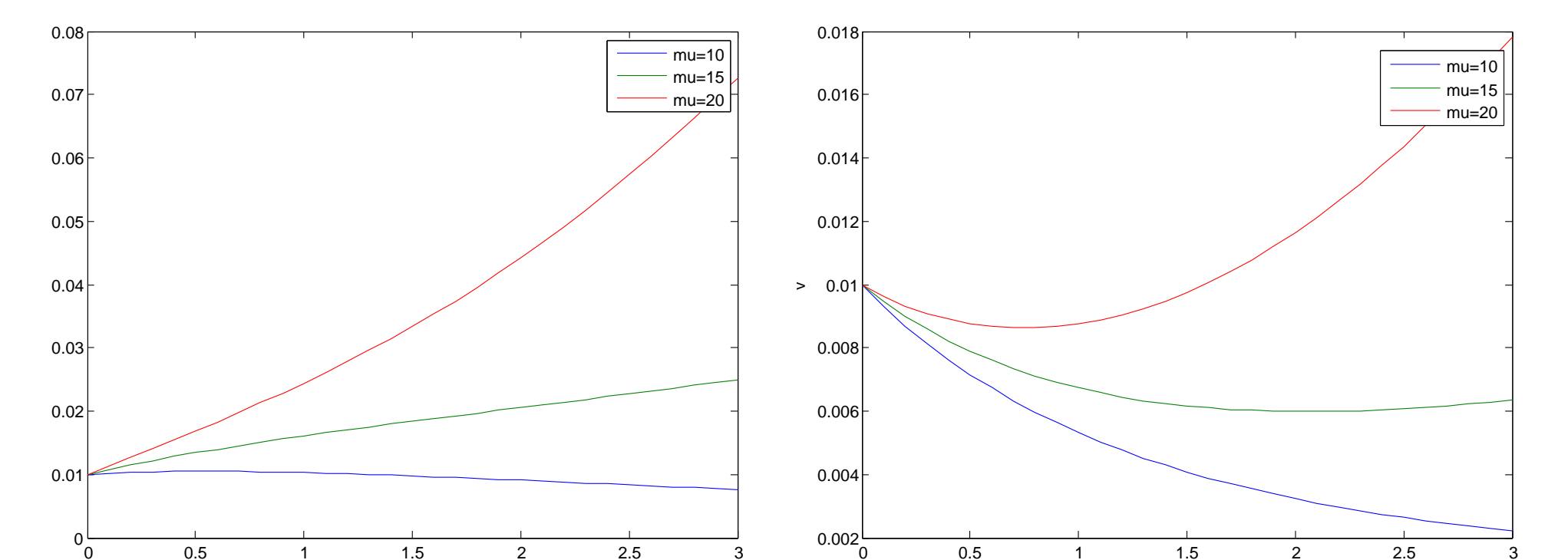


Figure 5. Evolution of the solution at the middle (left) and at the boundary (right) of the domain, for different values of μ .

Conclusion and Future Work

- ▶ A remarkable feature of our method is that we use an **implicit second order scheme** for the time discretisation, which improves its **accuracy and stability**, when compared with the available algorithms.
- ▶ To **reduce the computational complexity** of our method and improve its efficiency we have used an interpolation procedure which allows a **drastic reduction of matrix dimensions**, without a significant loss of accuracy.
- ▶ Our numerical results **confirm the theoretical predictions** and are in agreement with the expected behaviour of the solutions.
- ▶ As future work, we intend to analyse a **stochastic version** of the neural field equations, to take into account the **effect of noise**.

References

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