# Reduction and Lagrangian mechanics on Lie algebroids 

Patrícia Santos<br>CMUC, University of Coimbra<br>Engineering Institute of Coimbra

Mechanics and Lie algebroids, one day workshop September 11th, 2007

## Outline

(1) Reduction of Lie algebroids \& Poisson reduction
(2) Reduction of Lagrangian mechanics on Lie algebroids
(3) Hamel symbols and nonholonomic mechanics on Lie algebroids

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## (2) Reduction of Lagrangian mechanics on Lie algebroids

(3) Hamel symbols and nonholonomic mechanics on Lie algebroids

## Reduction of Lie algebroids

## Notation

$\checkmark(A, p, M),(\widehat{A}, \widehat{p}, \widehat{M})$ vector bundles
$\checkmark(\Pi, \pi):(A, p, M) \rightarrow(\widehat{A}, \widehat{p}, \widehat{M})$ epimorphism

Definition (Reduced Lie algebroid)
Iet $A$ and $\widehat{A}$ be $I$ ie algebroids with exterior derivatives $d_{A}$ and $d_{\lambda}$ respectively. The bundle $A$ is a reduced Lie algebroid of $A$ if $(\Pi, \pi)$ is a Lie algebroid homomorphism ${ }^{a}$, i.e. $d_{A} \circ \Pi^{*}=\Pi^{*} \circ d_{\hat{A}}$.
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## Definition (Section of an epimorphism)

A section of $\Pi$ is a set of maps between fibers

$$
S=\left\{\widehat{S_{x}^{x}}: \widehat{A}_{\widehat{x}} \rightarrow A_{x} \mid \pi(x)=\widehat{x}, x \in M\right\}
$$

such that $\Pi_{x} \circ S_{x}^{\widehat{x}}=\mathrm{id}_{\widehat{\mathrm{A}_{\hat{x}}}}$, for all $x \in M$.
(1) The set $S$ does not have to define a vector bundle morphism, but when it does $S$ is a section in the usual sense.
(2) A section $S$ defines a map $S: \Gamma(\widehat{A}) \rightarrow \Gamma(A)$ in the following way
$\mathbf{S}(\hat{v})(x)=S_{x}^{\hat{\hat{v}}}\left(\widehat{v_{\hat{x}}}\right)$,
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## Notes

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## Reduction theorem

$\checkmark\left(A, \rho,[\cdot, \cdot]_{A}\right)$ Lie algebroid over $M$
$\checkmark(\Pi, \pi):(A, p, M) \rightarrow(\widehat{A}, \widehat{p}, \widehat{M})$ surjective submersion
(1) $\exists(\widetilde{\Pi}, \pi):\left(A^{*}, \tau, M\right) \rightarrow\left(\widehat{A}^{*}, \widehat{\tau}, \widehat{M}\right)$ surjective submersion, such that $S=\left\{S_{x}^{\hat{x}}:=\left(\widetilde{\Pi}_{x}\right)^{*}: \widehat{A}_{\widehat{x}} \rightarrow A_{x} \mid \pi(x)=\widehat{x}, x \in M\right\}$ is a section of $\Pi$
(2) $\operatorname{Im} \mathbf{S}$ is a subalgebra of $\left(\Gamma(A),[\cdot, \cdot]_{A}\right)$
(3) $d_{A} \circ \Pi^{*} \circ \widetilde{\Pi}=\Pi^{*} \circ \widetilde{\Pi} \circ d_{A}$, where $\widetilde{\Pi}: \Omega^{\Pi}(A) \rightarrow \Omega(\widehat{A})$ is defined by
for all П-projectable $A$-k-form $\alpha$.

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\widetilde{\Pi}(\alpha)\left(\widehat{v}_{1}, \ldots, \widehat{v}_{k}\right)=\alpha\left(\mathbf{S}\left(\widehat{v}_{1}\right), \ldots, \mathbf{S}\left(\widehat{v}_{k}\right)\right)
$$

for all $\widetilde{\Pi}$-projectable $A$-k-form $\alpha$.

## Reduction theorem

Then, $\widehat{A}$ is a reduced Lie algebroid of $A$ with exterior derivative

$$
d_{\widehat{A}}:=\widetilde{\Pi} \circ d_{A} \circ \Pi^{*},
$$

and Lie algebroid structure $\left(\widehat{\rho},[\cdot, \cdot]_{\widehat{A}}\right)$ :

- $\widehat{\rho}=T \pi \circ \rho \circ S$
- $[\cdot, \cdot]_{\widehat{A}} \circ \pi=\Pi \circ[\mathbf{S}(\cdot), \mathbf{S}(\cdot)]_{A}$

An equivalent version of this theorem was proved by David Iglesias, Juan Carlos Marrero, David Martín de Diego, Eduardo Martínez and Edith Padrón, SIGMA 3 (2007).

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## Example: Poisson manifold

$\checkmark(M, \Lambda)$ Poisson manifold and $\left(T^{*} M, \Lambda^{\sharp},[\cdot, \cdot]_{T^{*} M}\right)$ associated Lie algebroid $\checkmark \pi: M \rightarrow \widehat{M}$ surjective submersion and $\sigma: \widehat{M} \rightarrow M$ is such that $\pi \circ \sigma=\mathrm{id}_{\widehat{\mathrm{M}}}$ $\checkmark \Pi=(T \sigma)^{*}: T^{*} M \rightarrow T^{*} \widehat{M}$ and $\widetilde{\Pi}=T \pi: T M \rightarrow T \widehat{M}$ are submersions over $\pi$
$\checkmark \operatorname{Im} \mathbf{S}$ is a Lie subalgebra of $\left(\Gamma\left(T^{*} M\right),[\cdot, \cdot]_{T^{*} M}\right)$ and the space of sections of $C_{\sigma}=\operatorname{Ker}(T \sigma)^{*}$ is an ideal of this algebra.

The reduced structure on $T^{*} \widehat{M}$ :

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The reduced structure on $T^{*} \widehat{M}$ :
(i) $\widehat{\rho}=T \pi \circ \Lambda^{\sharp} \circ S, \quad$ with $S_{x}^{\pi(x)}=\left(T_{x} \pi\right)^{*}: T_{\pi(x)}^{*} \widehat{M} \rightarrow T_{x}^{*} M$;
(ii) $[\cdot, \cdot]_{\widehat{A}} \circ \pi=\Pi \circ[\mathbf{S}(\cdot), \mathbf{S}(\cdot)]_{T^{*} M}, \quad$ with $\mathbf{S}=\pi^{*}: \Gamma\left(T^{*} \widehat{M}\right) \rightarrow \Gamma\left(T^{*} M\right)$.

## Poisson reduction

## Notation

$\checkmark\left(A, \rho,[\cdot, \cdot]_{A}\right)$ Lie algebroid over $M$
$\checkmark\left(A^{*},\{\cdot, \cdot\}_{A^{*}}\right)$ associated Lie co-algebroid

## Poisson reduction in the sense of Marsden-Ratiu '86

I et $C$ be a subbundle of $T A^{*}$ such that:
(1) it defines a surjective submersion $(\bar{\Pi}, \pi):\left(A^{*}, \tau, M\right) \rightarrow\left(A^{*}, \widehat{\tau}, \widehat{M}\right)$
(2) $\forall F, G \in C^{\infty}\left(A^{*}\right)$ such that $d F, d G \in C^{0}$, then $d\{F, G\}_{A^{*}} \in C^{0}$


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$\left.\widehat{A^{*}}, \widehat{\tau}, \widehat{M}\right)$ is endowed with a Poisson structure $\{\cdot, \cdot\}_{\wedge}$ that satisfies:
$\{\widehat{F} \circ \bar{\Pi}, \widehat{G} \circ \bar{\Pi}\} \quad=\{\widehat{F}, \widehat{G}\} \circ \bar{\Pi} . \quad \forall \widehat{F}, \widehat{G} \in C^{\infty}\left(\widehat{A^{*}}\right)$

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Theorem

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\left(A, \rho,[\cdot, \cdot]_{A}\right) \xrightarrow{L A R}\left(\widehat{A}, \widehat{\rho},[\cdot, \cdot]_{\widehat{A}}\right) \Longrightarrow\left(A^{*}, \Lambda_{A^{*}}\right) \xrightarrow{P R}\left(\widehat{A}^{*}, \Lambda_{\widehat{A}}\right)
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## Proposition



- $\widehat{\rho}=T \pi \circ \rho \circ \bar{S}$
- $\left.\overline{\mathbf{S}}(\Gamma . .]_{A}\right)=[\overline{\mathbf{S}}(.) \overline{\mathbf{\sigma}}(\cdot)]_{A}$
where $\bar{S}=\left\{\bar{S}_{x}^{\hat{x}}:=\left(\bar{\Pi}_{x}\right)^{*}: \widehat{A}_{\widehat{x}} \rightarrow A_{x} \mid \pi(x)=\widehat{x}, x \in M\right\}$.



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## Example: Lie group $G$

Let $G$ be finite dimensional Lie group with Lie algebra $\mathfrak{g}$.

- $(\Pi, \pi):(T G, p, G) \rightarrow(\mathfrak{g}, \widehat{p},\{\cdot\})$ is the canonical projection defined by the tangent representation $\Phi(g)=T L_{g}$ of the action of $G$ on itself by left translations, with $T G \equiv G \times \mathfrak{g}$ we have $\Pi(g, X)=X$ for all $X \in \mathfrak{g}$ and $g \in G$.
- With $T^{*} G \equiv G \times \mathfrak{g}^{*}$ the projection $\widetilde{\Pi}: T^{*} G \rightarrow \mathfrak{g}^{*}=T^{*} G / G$ defined by $\widetilde{\Pi}(g, \alpha)=\alpha$, for all $g \in G$ and $\alpha \in \mathfrak{g}^{*}$, is a surjective submersion over $\pi: G \rightarrow G / G=\{\cdot\}$.
By Marsden-Ratiu $g^{*}$ is endowed with a linear Poisson structure such that $\widetilde{\Pi}: T^{*} G \rightarrow \mathfrak{g}^{*}$ is a Poisson morphism. Then, we can prove that the conditions of the reduction theorem are satisfied and $\mathfrak{g}$, with its usual structure of Lie algebroid, is a reduced Lie algebroid of $T G$.


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(1) Reduction of Lie algebroids \& Poisson reduction
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(3) Hamel symbols and nonholonomic mechanics on Lie algebroids

## Reduction of a Lie algebroid prolongation

$\checkmark(\Pi, \pi):(A, p, M) \rightarrow(\widehat{A}, \widehat{p}, \widehat{M})$ submersion
$\checkmark \widehat{A}$ reduced Lie algebroid of $A$
$\mathcal{T} \Pi=(\Pi, \Pi, T \Pi): \mathcal{T} A \rightarrow \mathcal{T} \widehat{A}$ is a Lie algebroid surjective homomorphism over
$\mathcal{T} \widehat{A}$ is a reduced Lie algebroid of $\mathcal{T} A$.

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## Prolongation $\longleftrightarrow$ Reduction by symmetry

$\checkmark\left(A, \rho,[\cdot, \cdot]_{A}\right)$ Lie algebroid over $M$
$\checkmark \Phi: G \rightarrow \operatorname{Aut}(A)$ Lie algebroid representation of $G$ on $A$
$\checkmark \Phi$ e $\Phi^{c}$ define proper and free actions of $G$ on $A$ and $A^{*}$
$A$ satisfies the reduction theorem conditions, and then $A / G$ is a reduced Lie algebroid of $A$ over $M / G$.

## $\mathcal{T}(A / G) \simeq(\mathcal{T} A) / G$ is a reduced Lie algebroid of $\mathcal{T} A .{ }^{a}$

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$\mathcal{T}(A / G) \simeq(\mathcal{T} A) / G$ is a reduced Lie algebroid of $\mathcal{T} A .^{a}$
$\frac{{ }^{a} \text { Using a different approach, this result was also proven by Manuel de León, Juan Carlos Marrero and }}{\text { Eduardo Martínez, J. Phys. A: Math Gen. } 38 \text { (2005). }}$.

## Example: Principal fiber bundle $P(M, G)$

Consider a principal fibre bundle $P(M, G)$ and the associated gauge algebroid (TP/G, $p, M$ ).

- The canonical projection $\Pi: T P \rightarrow T P / G$ is a homomorphism of Lie algebroids over $\pi: P \rightarrow M$ that defines the homomorphism of Lie algebroids $\mathcal{T} \Pi: \mathcal{T}(T P) \rightarrow \mathcal{T}(T P / G)$ over $\Pi$. Note that $\mathcal{T}(T P) \equiv T(T P)$.
- Let $\phi$ be the (right) action of the Lie group $G$ on $P$. Then, $\Phi(g):=T \phi_{g}$ defines a Lie algebroid representation of $G$ on $T P$.

We can prove that $\mathcal{T} \Phi=(\Phi, \Phi, T \Phi)$ is a Lie algebroid representation of $G$ on $\mathcal{I}(T P)$ and then $\mathcal{T}(T P) / G \cong \mathcal{T}(T P / G)$, that is, $T(T P) / G \cong \mathcal{T}(T P / G)$.

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## Reduction of Lagrangian mechanics on Lie algebroids

## Notation

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$\checkmark(\Pi, \pi):(A, p, M) \rightarrow(\widehat{A}, \widehat{p}, \widehat{M})$ submersion

## Lemma

$W / e$ have $\mathcal{T}\|\circ S=\hat{S} \circ T\|$ and $\mathcal{T} \Pi \circ \Delta=\Delta \circ \Pi$, where $\hat{S}$ and $\Delta$ are, respectively, the vertical endomorphism and the Liouville section of $\mathcal{T} A$.

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## Lemma

We have $\mathcal{T} \Pi \circ S=\widehat{S} \circ \mathcal{T} \Pi$ and $\mathcal{T} \Pi \circ \Delta=\widehat{\Delta} \circ \Pi$, where $\widehat{S}$ and $\widehat{\Delta}$ are, respectively, the vertical endomorphism and the Liouville section of $\mathcal{T} \widehat{A}$.

## Theorem

Let $L \in C^{\infty}(A)$ be a $\Pi$-invariant Lagrangian of a dynamics Lagrangian system on $A$, i.e. there exists $l \in C^{\infty}(\widehat{A})$ such that $L=l \circ \Pi$. Then:

- $\widehat{E}_{l} \circ \Pi=E_{L}$
- $(\mathcal{T} \Pi)^{*} \widehat{\theta}_{l}=\theta_{L} \Longrightarrow(\mathcal{T} \Pi)^{*} \widehat{\omega}_{l}=\omega_{L}$


## When $L$ is regular, then we have:

- $l$ is regular

Therefore, the Lagrangian dynamics on $A$ induced by a regular $\Pi$-invariant Lagrangian $L=l \circ \Pi$ reduces to the dynamics on $\widehat{A}$ given by $l$.

## Theorem

Let $L \in C^{\infty}(A)$ be a $\Pi$-invariant Lagrangian of a dynamics Lagrangian system on $A$, i.e. there exists $l \in C^{\infty}(\widehat{A})$ such that $L=l \circ \Pi$. Then:

- $\widehat{E}_{l} \circ \Pi=E_{L}$
- $(\mathcal{T} \Pi)^{*} \widehat{\theta}_{l}=\theta_{L} \Longrightarrow(\mathcal{T} \Pi)^{*} \widehat{\omega}_{l}=\omega_{L}$

When $L$ is regular, then we have:

- $l$ is regular
- $\mathcal{T} \Pi \circ X_{L}=\widehat{X}_{l} \circ \Pi$

Therefore, the Lagrangian dynamics on $A$ induced by a regular $\Pi$-invariant Lagrangian $L=l \circ \Pi$ reduces to the dynamics on $\widehat{A}$ given by $l$.

## Outline

## (1) Reduction of Lie algebroids \& Poisson reduction

(2) Reduction of Lagrangian mechanics on Lie algebroids
(3) Hamel symbols and nonholonomic mechanics on Lie algebroids

## Adapted coordinates to nonholonomic restrictions

## Notation

$\checkmark\left(A, \rho,[\cdot, \cdot]_{A}\right)$ Lie algebroid over $M$
$\checkmark k$ nonholonomic linear restritions on $A$

$$
\phi_{a}(q, \mathbf{v})=\widehat{\Phi_{a}}(q, \mathbf{v})=\phi_{a \beta}(q) \mathbf{v}^{\beta},
$$

where $\widehat{\Phi_{a}}$ is a linear function defined by the $A$-1-form $\Phi_{a}$

Adapted coordinates
$\left\{\left(q^{i} \mathbf{w}^{\alpha}\right) \mid i=1 \ldots,{ }^{n}, \alpha=1, \ldots, s\right\}$ local coordinates on $A$

- $\mathrm{w}^{\alpha}=\Phi_{\alpha \beta} \mathrm{v}^{\beta}$ the last $k$ coordinates coincide with the restrictions $\phi_{a}$, i.e.


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- $\left\{\left(q^{i}, \mathbf{w}^{\alpha}\right) \mid i=1, \ldots, n, \alpha=1, \ldots, s\right\}$ local coordinates on $A$
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$$
\begin{aligned}
\mathbf{w}^{I} & =\Phi_{I \beta} \mathbf{v}^{\beta}, \quad \forall I=1, \ldots,(s-k) \\
\mathbf{w}^{s-k+a} & =\phi_{a}, \quad \forall a=1, \ldots, k
\end{aligned}
$$

Associated with the new coordinates $A$, we consider on the prolongation of $A$ the following basis of local sections:

$$
\mathcal{X}^{\prime}{ }_{\alpha}(a)=\left(a, f_{\alpha}(p(a)), X_{\alpha}(a)\right), \quad \mathcal{V}^{\prime}{ }_{\alpha}(a)=\left(a, 0,\left.\frac{\partial}{\partial \mathbf{w}^{\alpha}}\right|_{a}\right),
$$

where $X_{\alpha}=\left.\Psi_{\beta \alpha} \rho^{i}{ }_{\beta} \partial_{q^{i}}\right|_{\mathbf{w}}$, for all $\alpha=1, \ldots, r$, where $\Psi_{\alpha \beta} \Phi_{\beta \gamma}=\delta_{\alpha \gamma}$.

## Lie algebroid structure on $\mathcal{T} A$

$$
\begin{array}{rlrl}
{\left[\mathcal{X}^{\prime}{ }_{\alpha}, \mathcal{X}^{\prime}{ }_{\beta}\right]_{\mathcal{T} A}} & =\gamma_{\alpha \beta}^{\epsilon} \mathcal{X}_{\epsilon}^{\prime}, & {\left[\mathcal{X}^{\prime}{ }_{\alpha}, \mathcal{V}_{\beta}^{\prime}\right]_{\mathcal{T A}}=0, \quad\left[\mathcal{V}_{\alpha}^{\prime}, \mathcal{V}^{\prime}{ }_{\beta}\right]_{\mathcal{T A}}=0,} \\
\rho_{\mathcal{T A}}\left(\mathcal{X}^{\prime}{ }_{\alpha}\right) & =X_{\alpha}, \quad \rho_{\mathcal{T A}}\left(\mathcal{V}^{\prime}{ }_{\alpha}\right)=\frac{\partial}{\partial \mathbf{w}^{\alpha}},
\end{array}
$$

where $\left[f_{\alpha}, f_{\beta}\right]_{A}=\gamma_{\alpha \beta}^{\epsilon} f_{\epsilon}$.

## Euler-Lagrange equations in adapted coordinates

Let $L \in C^{\infty}(A)$ be a regular Lagrangian of a dynamical system on the Lie algebroid $A$ with a non-conservative force $\mathcal{Q}$.

## Euler-Lagrange generalized equations:


with:

- $\dot{q}^{i}=W^{\alpha} \Psi_{\beta \alpha} \rho^{i} \beta$, where $\Psi=\Phi^{-1}$
- $\Upsilon_{\alpha}$ is the $\alpha$-component of the nonconservative force $\mathcal{Q}$, in the new coordinates

Mramel symbols

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Euler-Lagrange generalized equations:

$$
\frac{d}{d t}\left(\frac{\partial L}{\partial \mathbf{w}^{\alpha}}\right)=\Psi_{\beta \alpha} \rho^{i}{ }_{\beta} \frac{\partial L}{\partial q^{i}}+\mathbf{w}^{\epsilon} \gamma_{\epsilon \alpha}^{\beta} \frac{\partial L}{\partial \mathbf{w}^{\beta}}+\Upsilon_{\alpha},
$$

with:

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- $\gamma_{\epsilon \alpha}^{\beta}$ Hamel symbols


## Gauge algebroid $T G / G$ of a Lie group $G$

Let $G$ be a Lie group and $e$ its neutral element.

- Using the map $T L_{g^{-1}}: T G \rightarrow G \times T_{e} G$ given by $T L_{g^{-1}}(g, \dot{g})=(g, \xi)$ we can identify $T G$ with $G \times T_{e} G$.
- If $\left(\xi^{I}\right)$, for $I=1, \ldots, \operatorname{dim} G$, is the set of coordinates of $\xi \in T_{e} G$ with respect to a basis $\left\{e_{I}\right\}$ of $T_{e} G$, then we can define a set of quasi-velocities ( $\xi^{I}$ ) on $T G$ by $\xi^{I} e_{I}=\xi=T_{g} L_{g^{-1}}(\dot{g})$. If $g$ is a point in $G$ of local coordinates $\left(g^{I}\right)$, then $\left(g^{I}, \xi^{I}\right)$ defines a set of quasi-coordinates in $T G$.

A regular $G$-invariant Lagrangian $\mathcal{L} \in C^{\infty}(T G)$ in quasi-coordinates is given by $\mathcal{L}(g, \dot{g})=l(\xi)$, where $l$ is a function on the Lie algebra $\mathfrak{g}=T_{e} G$. The Euler-Lagrange equations of the gauge algebroid $T G / G \equiv \mathfrak{g}$ are given by:

where $c_{J I}{ }^{K}$ are the structure constants of the Lie algebra $g$ of the Lie group $G$ with resnect to the hasic $\left\{e_{i}\right\}$ of $g$.

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$$
\frac{d}{d t}\left(\frac{\partial l}{\partial \xi^{I}}\right)=\xi^{J} c_{J I}^{K} \frac{\partial l}{\partial \xi^{K}},
$$

where $c_{J I}{ }^{K}$ are the structure constants of the Lie algebra $\mathfrak{g}$ of the Lie group $G$ with respect to the basis $\left\{e_{I}\right\}$ of $\mathfrak{g}$.


[^0]:    aUsing a different appronch, this result was also proven by Mamel de I cón, Juan Camlos Marrero and Eduardo Martínez, J. Phys. A: Math Gen. 38 (2005).

[^1]:    ${ }^{a}$ Using a different approach, this result was also proven by Manuel de León, Juan Carlos Marrero and Eduardo Martínez. J. Phys. A: Math Gen. 38 (2005).

