

Gaugino Condensation and Generation of Supersymmetric 3-Form Flux

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Study: gaugino condensation in
 $\mathcal{N} = 1^*$ $SU(N)$ gauge theory, large N , large
 $g_{YM}^2 N$

AdS/CFT correspondence \rightarrow Polchinski-Strassler
solution of type IIB supergravity

Dynamical effect \rightarrow growth of 3-form flux?

Show: Hodge type $(3,0)$, IASD \rightarrow supersymmetry?

$\mathcal{N} = 1^*$: mass deformation of $\mathcal{N} = 4$ $SU(N)$

$$W = m_{ij} \phi_{(i} \phi_{j)}, \quad m_{ij} = m, \quad i = 1, 2, 3$$

Confining phase: $\langle \bar{\lambda} \bar{\lambda} \rangle = m^3$ at large $g_{YM}^2 N$

Gaugino condensation: order m^3 effect in dual
supergravity solution

Dual supergravity solution describing RG-flow:
warped solution with 3-form flux

$$ds^2 = Z^{-1/2} \eta_{\mu\nu} dx^{\mu\nu} + Z^{1/2} g_{mn} dx^m dx^n ,$$

$$G_3 = F_3 - \tau H_3 , \quad F_3 = dC_2 , \quad H_3 = dB_2$$

$$G_3 = G_3(x^m) , \quad \tau = \tau(x^m)$$

Exact supergravity solution not known. Iterative expansion in powers of m

Polchinski-Strassler 2000: order m , $\tau = \tau_0$

Freedman/Minahan 2000: order m^2 , $d\tau \neq 0$

Here: order m^3

Radial coordinate \rightarrow energy scale in dual FT

$$r^2 = x^m x^m = \bar{z}^i z^i , \quad i = 1, 2, 3$$

Large $r \rightarrow$ UV regime

Mass deformation \rightarrow 2-form-potential

$$A_2 = C_2 - \tau_0 B_2 \quad (G_3 = dA_2)$$

$W = m \Phi_i^2 \rightarrow$ dimension $\Delta = 3$

$$A_2 = m r^{\Delta-4} \varepsilon_{i\bar{j}\bar{k}} \frac{z^i}{r} d\left(\frac{\bar{z}^j}{r}\right) \wedge d\left(\frac{\bar{z}^k}{r}\right)$$

Polarisation tensor of type $(1, 2)$

3-form flux $G_3 = dA_2$ has $(1, 2)_P, \dots$, but NO
 $(3, 0)$ -piece

Gaugino condensate (chiral primary)

$$\langle \bar{\lambda} \bar{\lambda} + \dots \rangle = m^3, \quad \Delta = 3$$

Singlet of $SU(3) \subset SU(4)_R$ of $\mathcal{N} = 4$ SYM

2-form potential

$$A_2^{(2,0)} = m^3 r^{-\Delta} \varepsilon_{ijk} f \left(\frac{z}{r}, \frac{\bar{z}}{r} \right) \frac{z^i}{r} d \left(\frac{z^j}{r} \right) \wedge d \left(\frac{z^k}{r} \right)$$

Polarisation tensor $(3,0)$, singlet of $SU(3)$

Determine f : type IIB EOM for $G_3 = dA_2$

Schematically

$$d [Z^{-1} (\star - i) G_3] = m d\tau \wedge d\bar{z}^i \wedge dz^j \wedge dz^k \varepsilon_{ijk}$$

Homogeneous and inhomogeneous solutions at
order m^3

Homogeneous solution: $(\star - i)G_3 = 0$

→ ISD

→ G_3 does NOT contain $(3,0)$ -piece

$f = c = \text{constant}$, undetermined at large r ,
fixed by physics in IR

Inhomogeneous solution: uniquely fixed

→ G_3 HAS $(3,0)$ -piece, $G^{(3,0)} \propto m \partial \tau$

CAN be written as $G^{(3,0)} = \partial A_2^{(2,0)}$,

$$A_2^{(2,0)} = m^3 r^{-\Delta} \varepsilon_{ijk} f \frac{z^i}{r} d\left(\frac{z^j}{r}\right) \wedge d\left(\frac{z^j}{r}\right)$$

$$f = \left(\frac{\bar{z}^2}{r^2}\right)^2$$

$m^3 : G^{(3,0)} \rightarrow \varepsilon_{ijk}$ gaugino condensate vev

Supersymmetry:

Gaugino condensate preserves $\mathcal{N} = 1$ supersymmetry. IASD $G^{(3,0)} \rightarrow$ supersymmetry?

Graña/Polchinski 2000,

Giddings, Kachru, Polchinski 2001:

supersymmetric G_3 -flux in type IIB is ISD, of type $(2,1)_P$.

Type B spinor ansatz:

one globally defined six-dimensional Weyl spinor $\eta \rightarrow$ transverse six-dimensional space has SU(3) group structure (J, Ω)

More general spinor ansatz: Dall'Agata 2004

Type D: $\eta, \chi \rightarrow$ SU(2) structure (J^i, w_m)

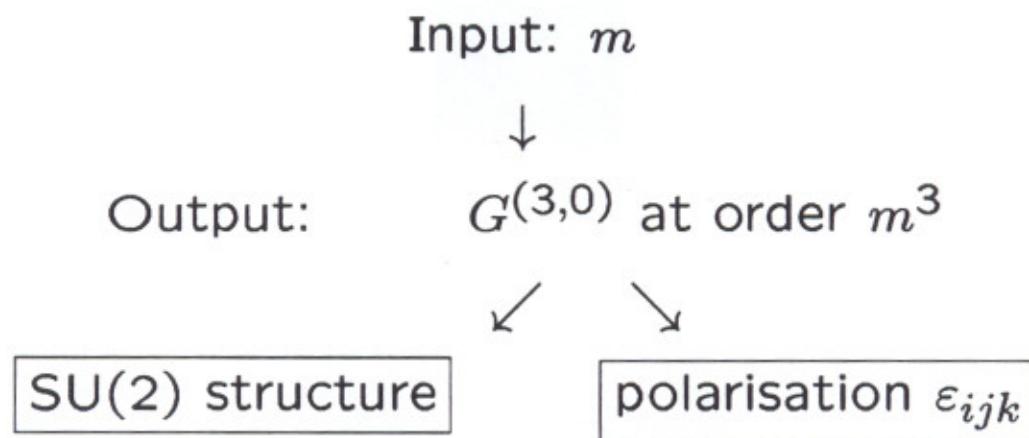
$$w_m = \chi^T \gamma_m \eta$$

Result:

PS background at order m^3 with IASD flux
 $G^{(3,0)}$ is supersymmetric, SU(2) structure

Type D ansatz: manifestation of dielectric na-
ture of solution Pilch and Warner 2004
Projector for D3 (UV) and D5/NS5 (IR)

Summarising:



SU(2) structure helpful in finding the exact type IIB solution associated to the $\mathcal{N} = 1^*$ gauge theory.